

Measuring black hole spin in AGN: a few recent X-ray driven results

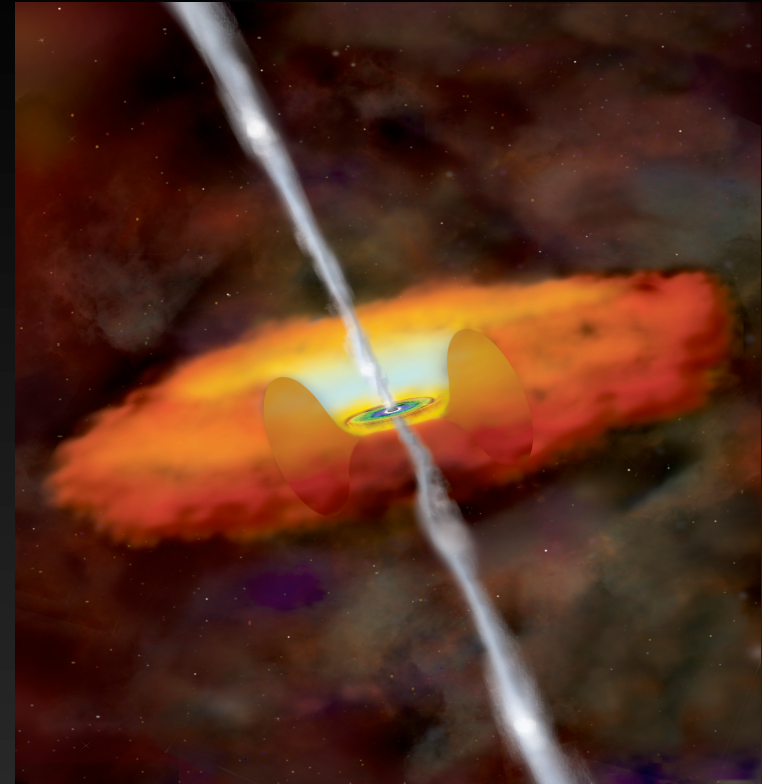
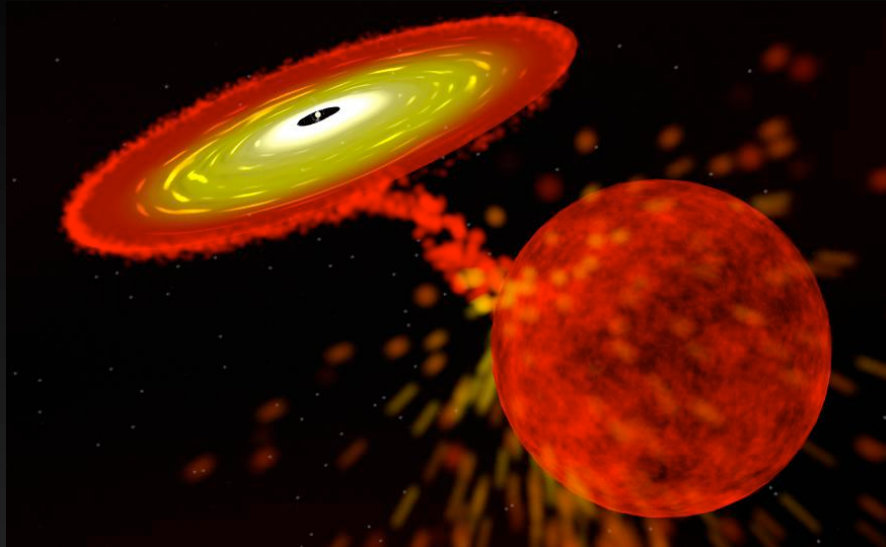
Giovanni Miniutti



Probing strong gravity around BH – Prague – Feb 2010



Accreting BHs

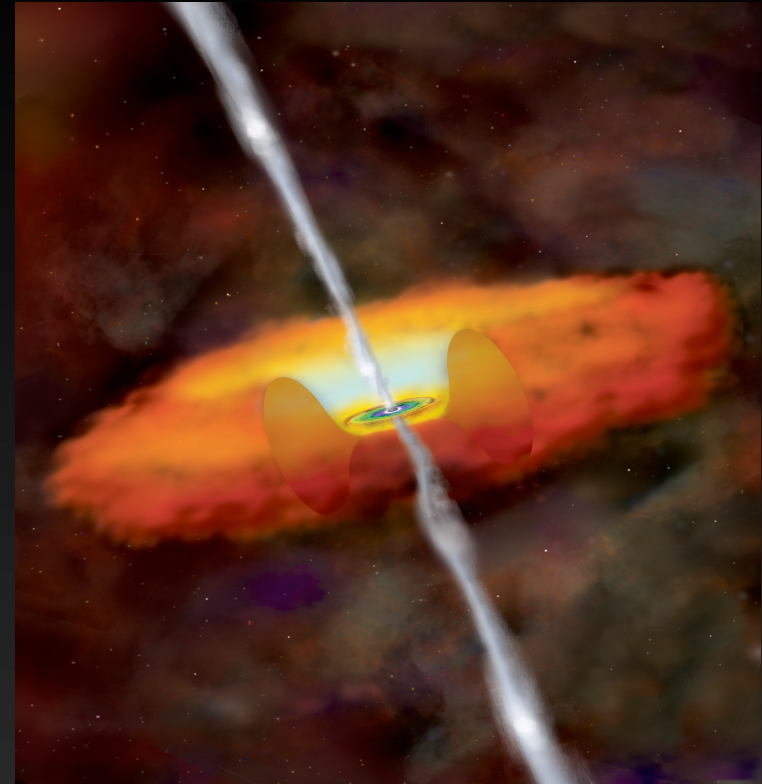
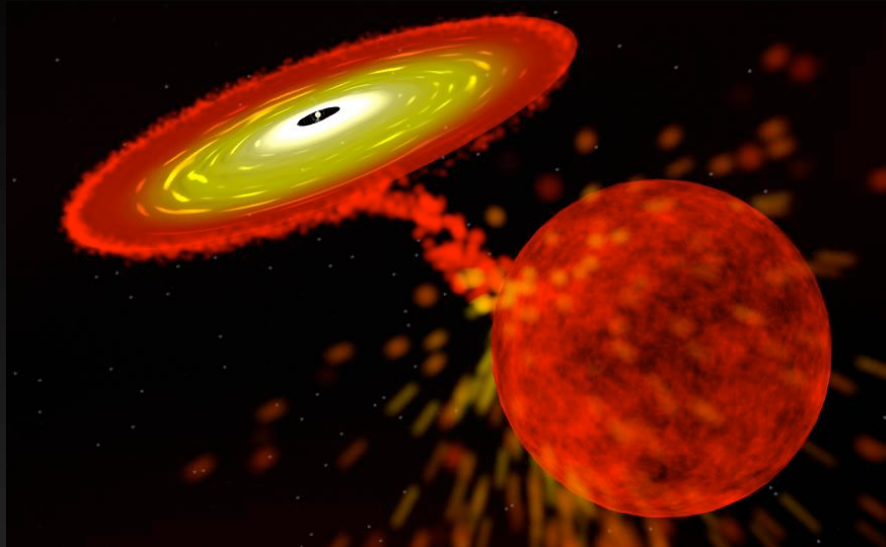


The two flavours of accreting BH:

- stellar mass BHs scattered in galaxies (X-ray binaries)
- supermassive BHs in the center of galaxies (AGN and quasars)



Accreting BHs

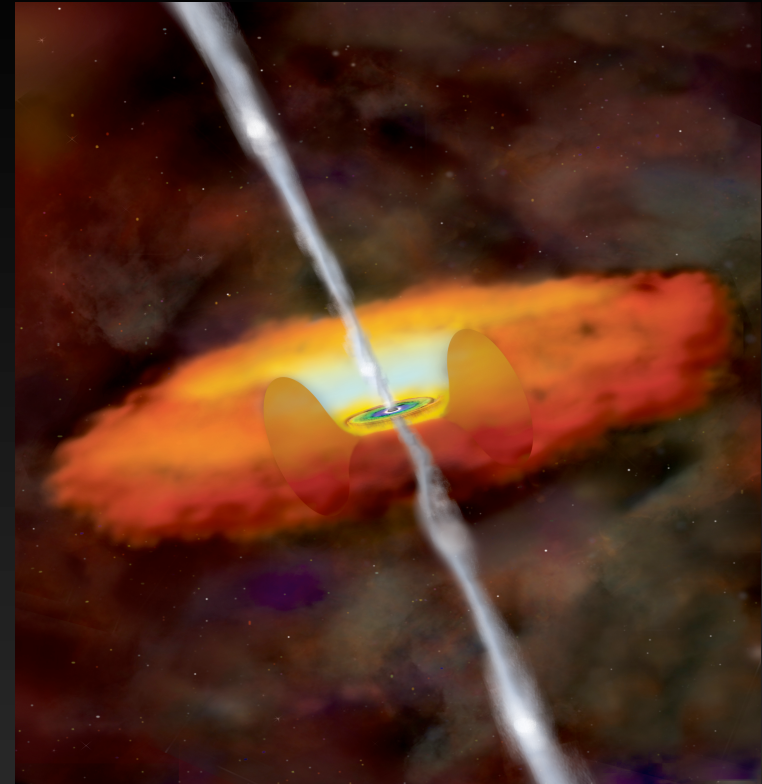
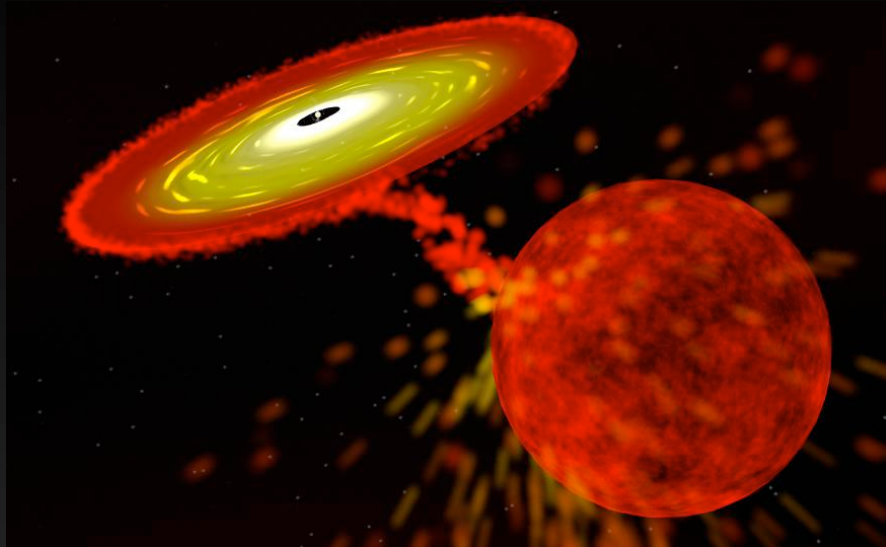


Why caring about the BH spin ?:

stellar mass BHs dynamics of BH formation in supernovae
supermassive BHs necessary condition for jets ?
relative importance of mergers and accretion
accretion modes (coherent vs chaotic)



Accreting BHs



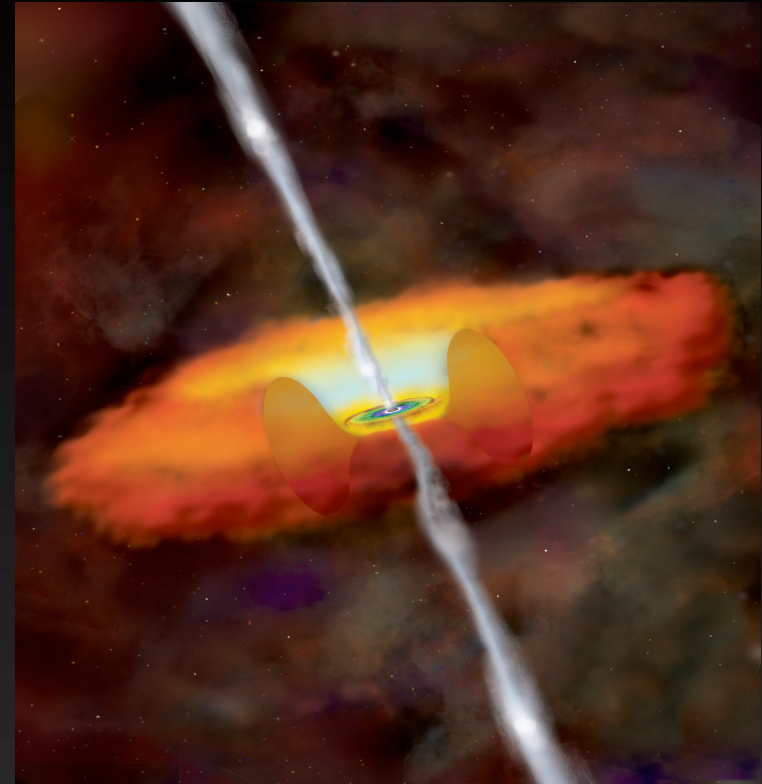
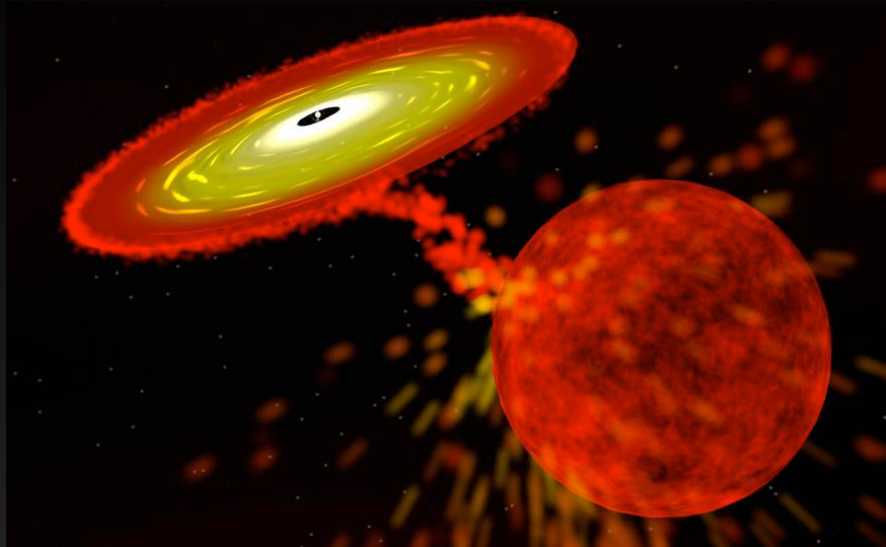
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stellar mass BHs thermal continuum, Fe K lines, QPO

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Accreting BHs



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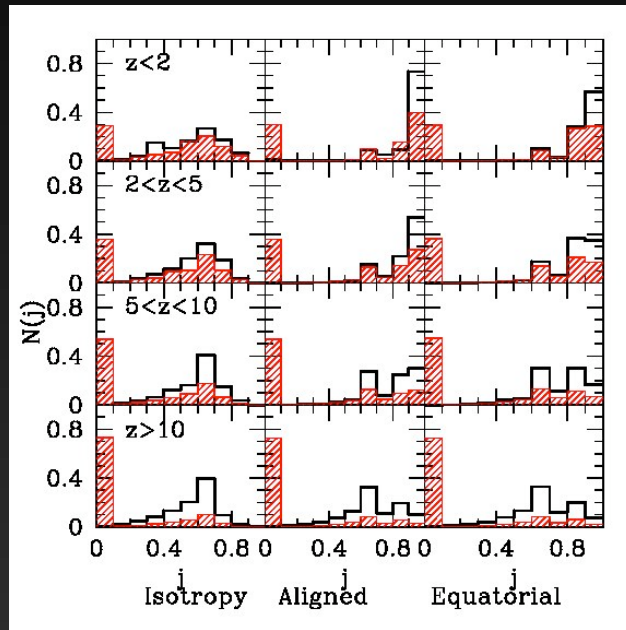
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supermassive BHs ~~thermal continuum~~, ~~Fe K lines~~, ~~QPO~~

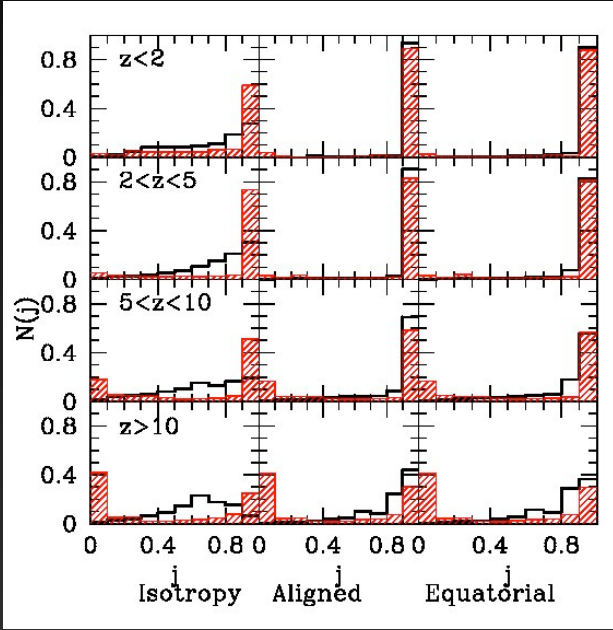


BH spin in AGN

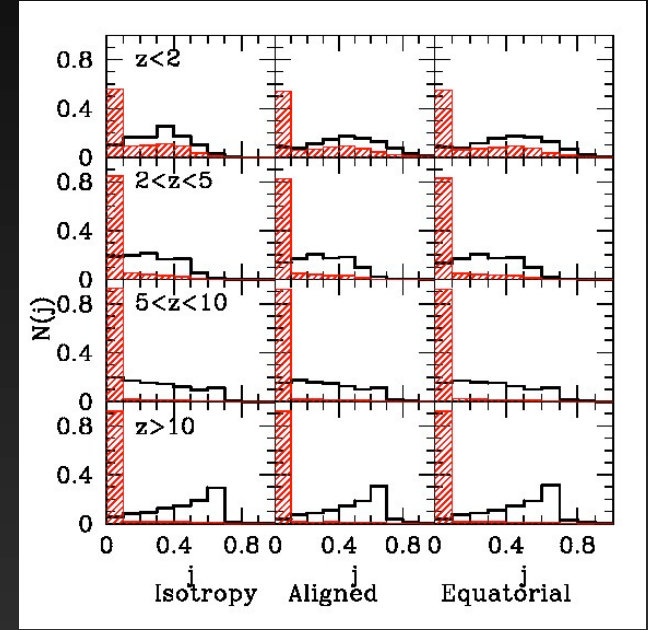
mergers only
 $a \sim 0.7$



mergers + coherent
high a



mergers + chaotic
low a

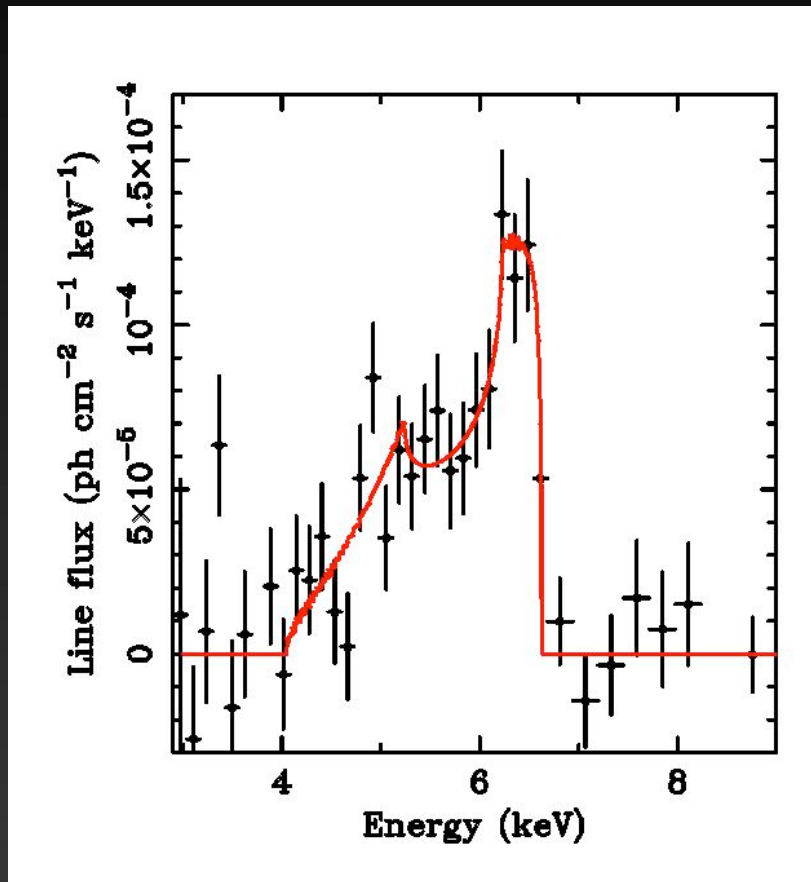


Berti & Volonteri 08



The usual suspect: MCG-6-30-15

First clear **detection of relativistic Fe K line** (Tanaka et al 95) and first evidences for a **rapidly spinning Kerr BH** (Iwasawa et al 96, 99)

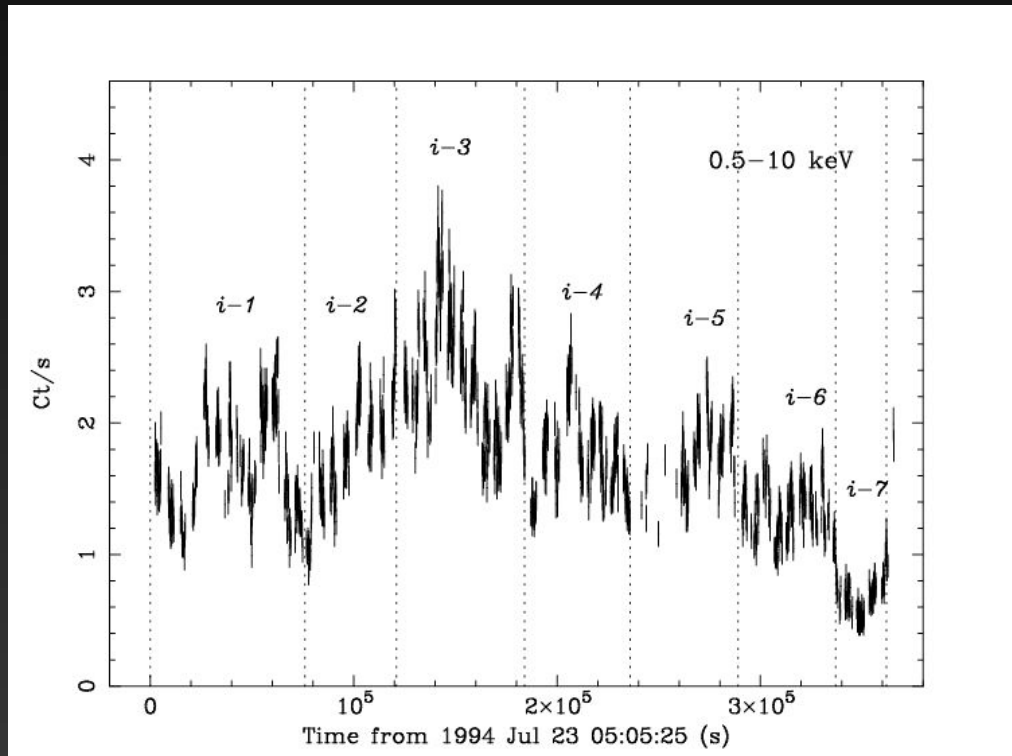


Tanaka et al 95

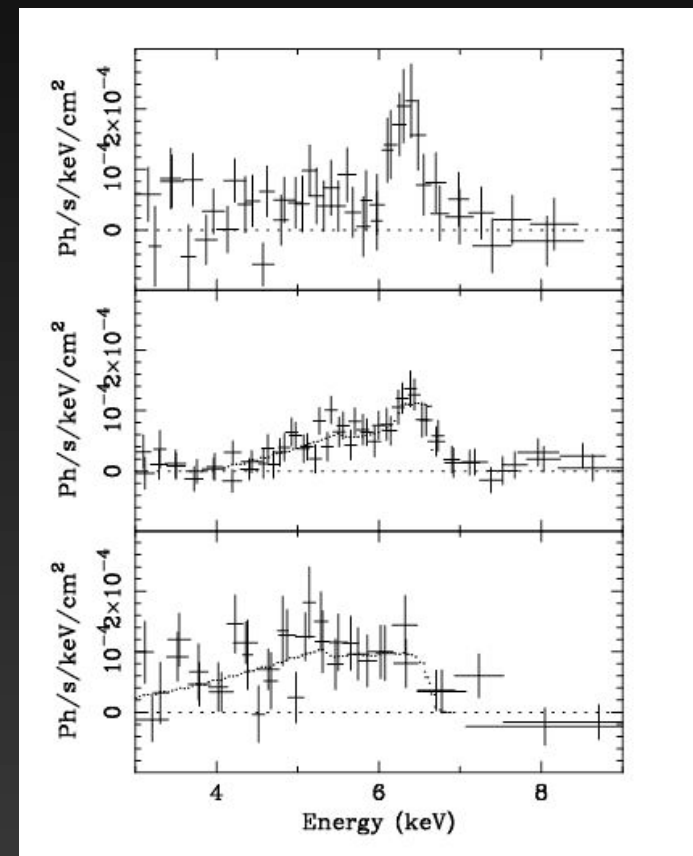


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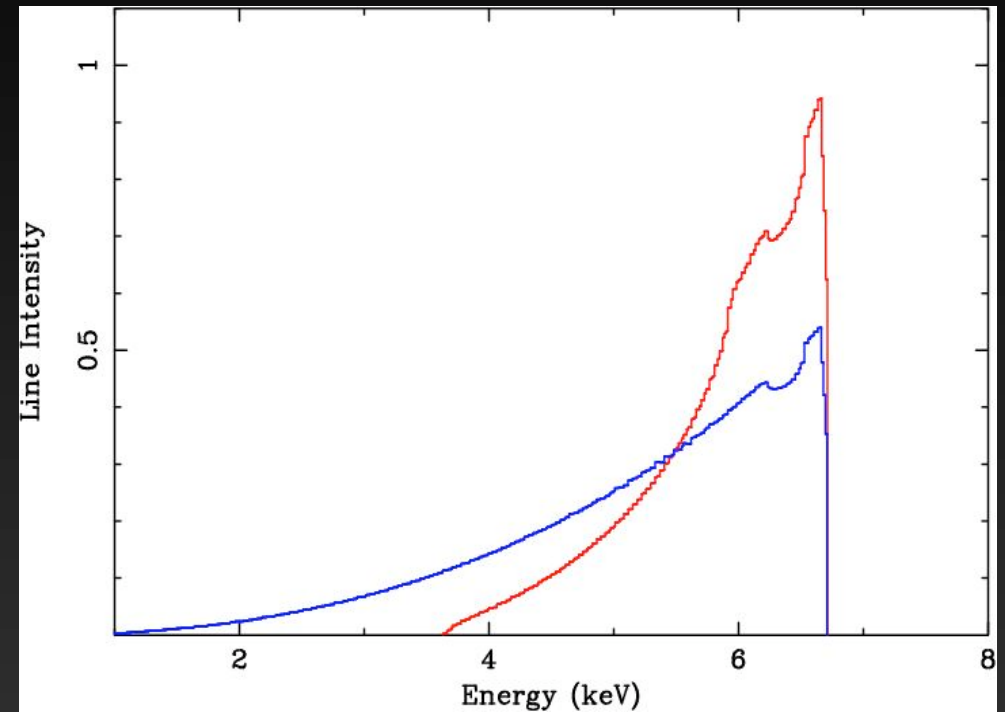
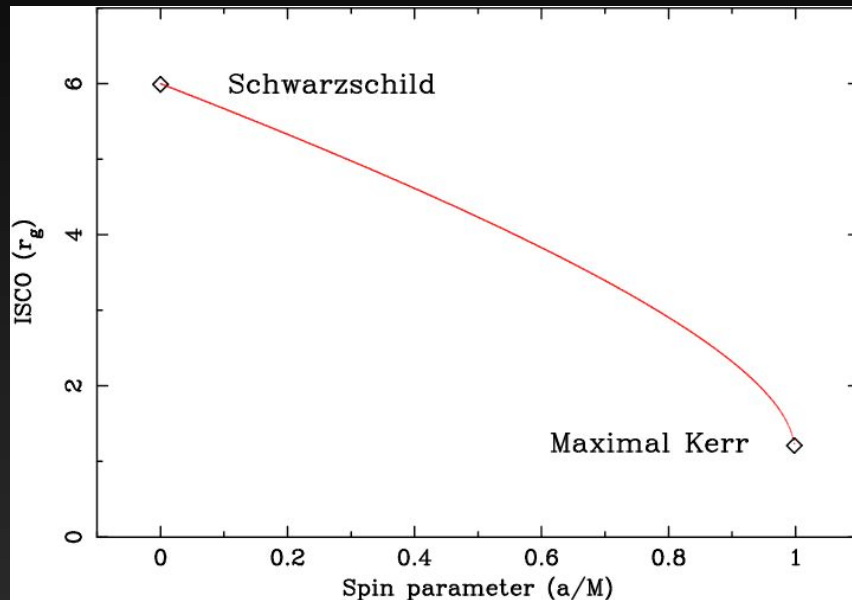
Iwasawa et al 96





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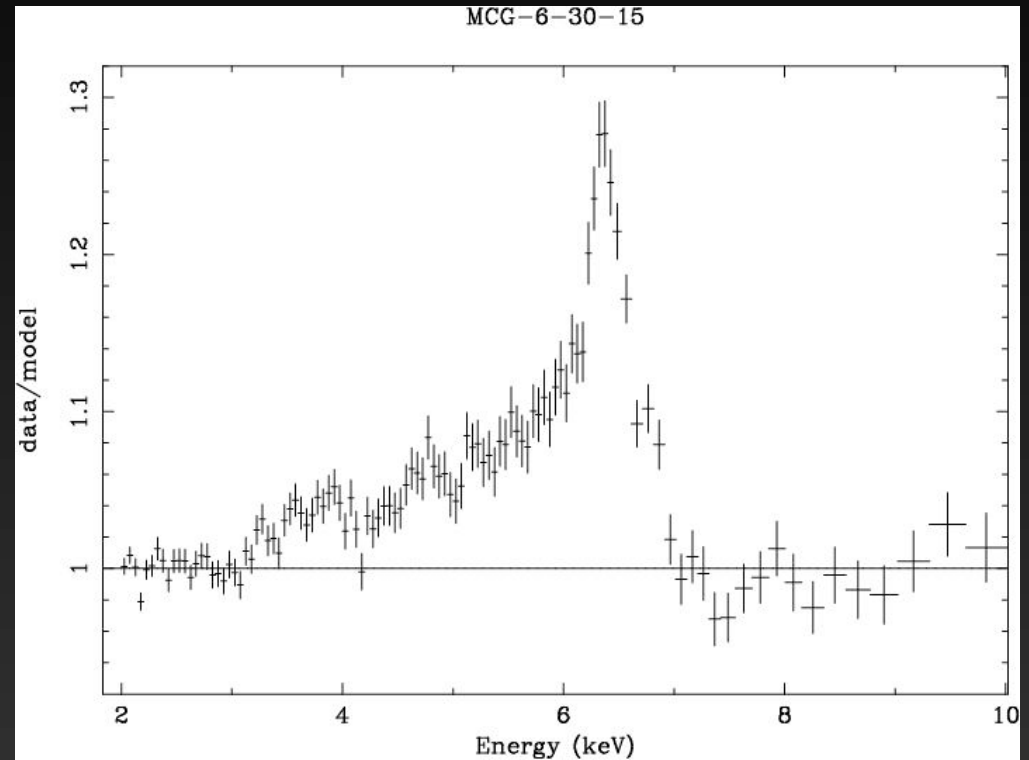
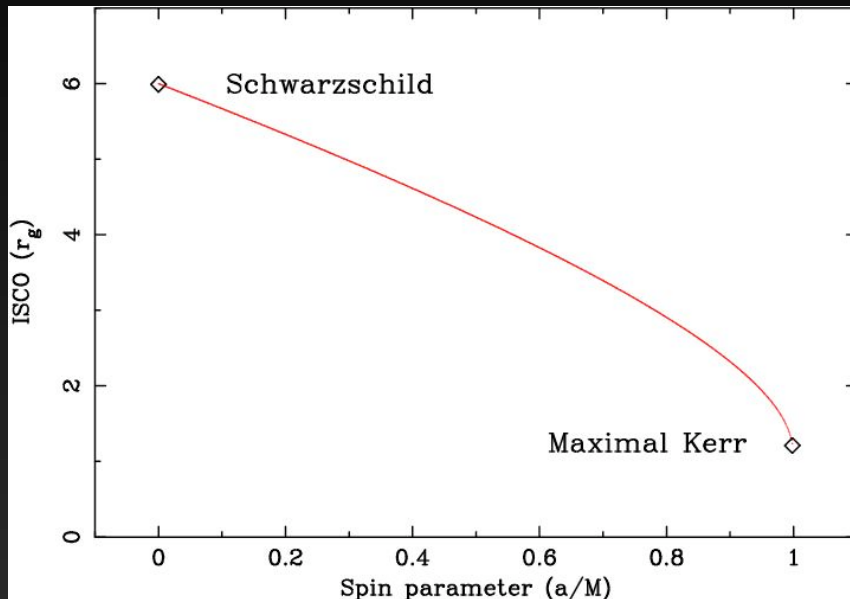
BH spin measurements rely on the id. $\text{ISCO} \approx R_{\text{in}}$





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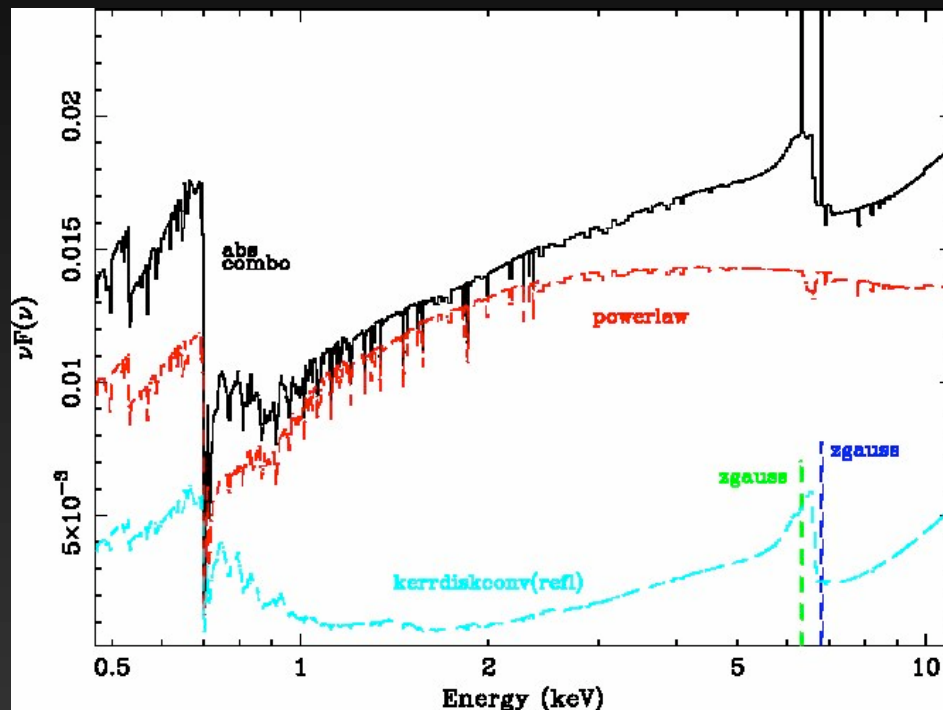
Early results in MCG-6 indicate that $R_{in} < 2 r_g$
which translates into a BH spin of $a > 0.94$

Fabian et al 02



The usual suspect: MCG-6-30-15

More recently, a detailed analysis by Brenneman & Reynolds (07) yields to a much more accurate spin measurement of $a \geq 0.987$

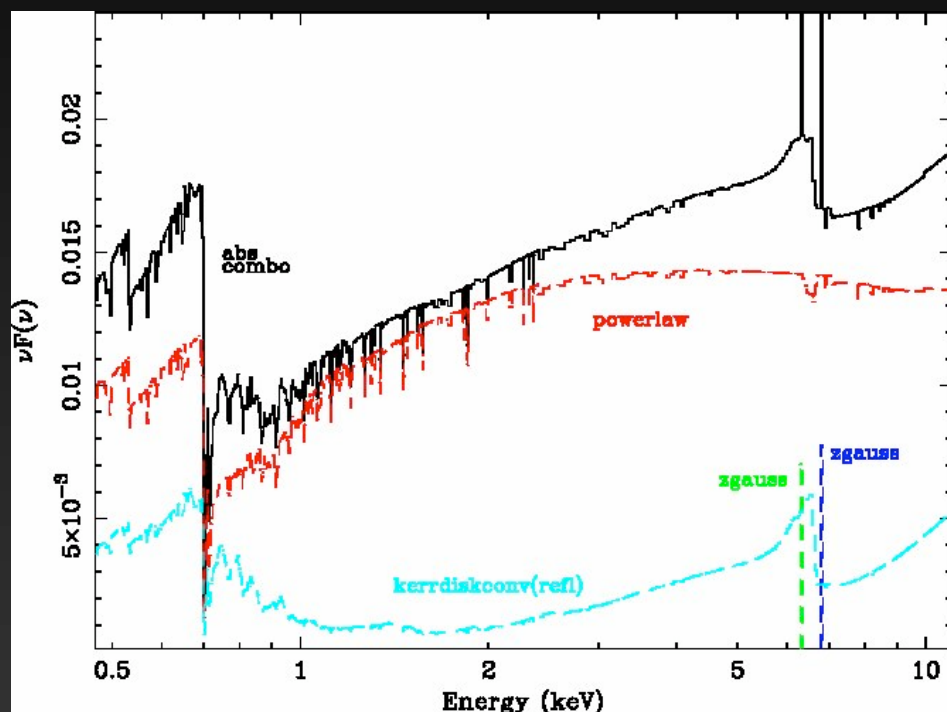


Brenneman & Reynolds 07
but see Miller et al. 09

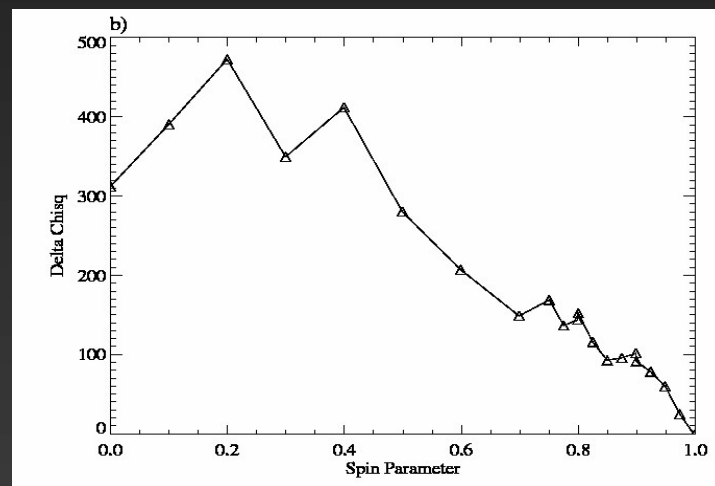
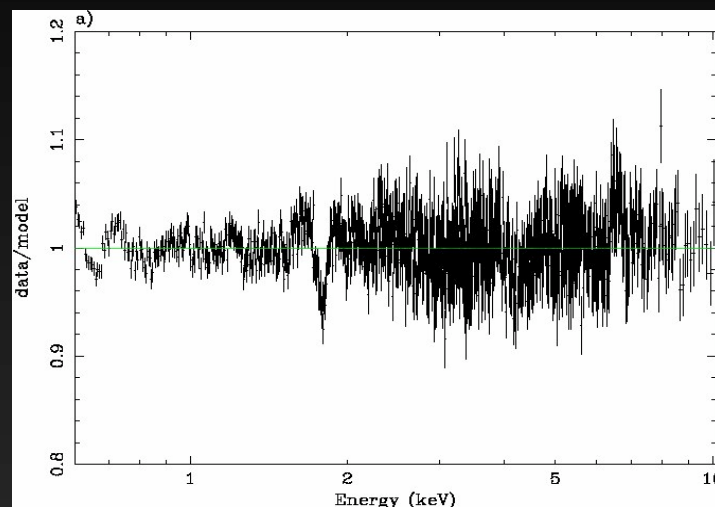


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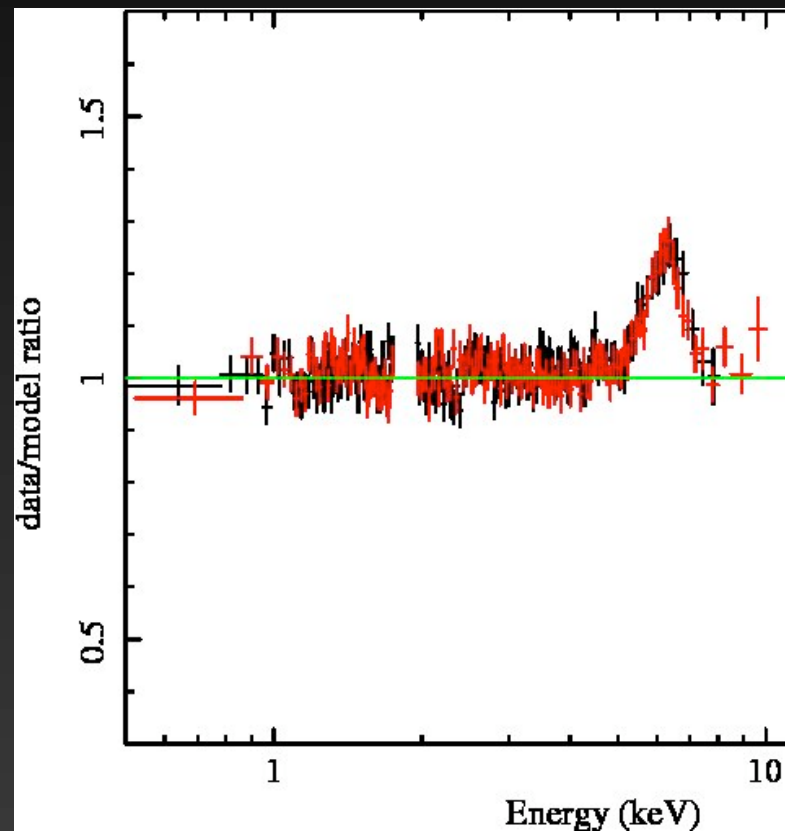




Swift J2127.4+5654 with Suzaku

This is a **NLS1** galaxy detected above 20 keV with INTEGRAL/IBIS (rare)

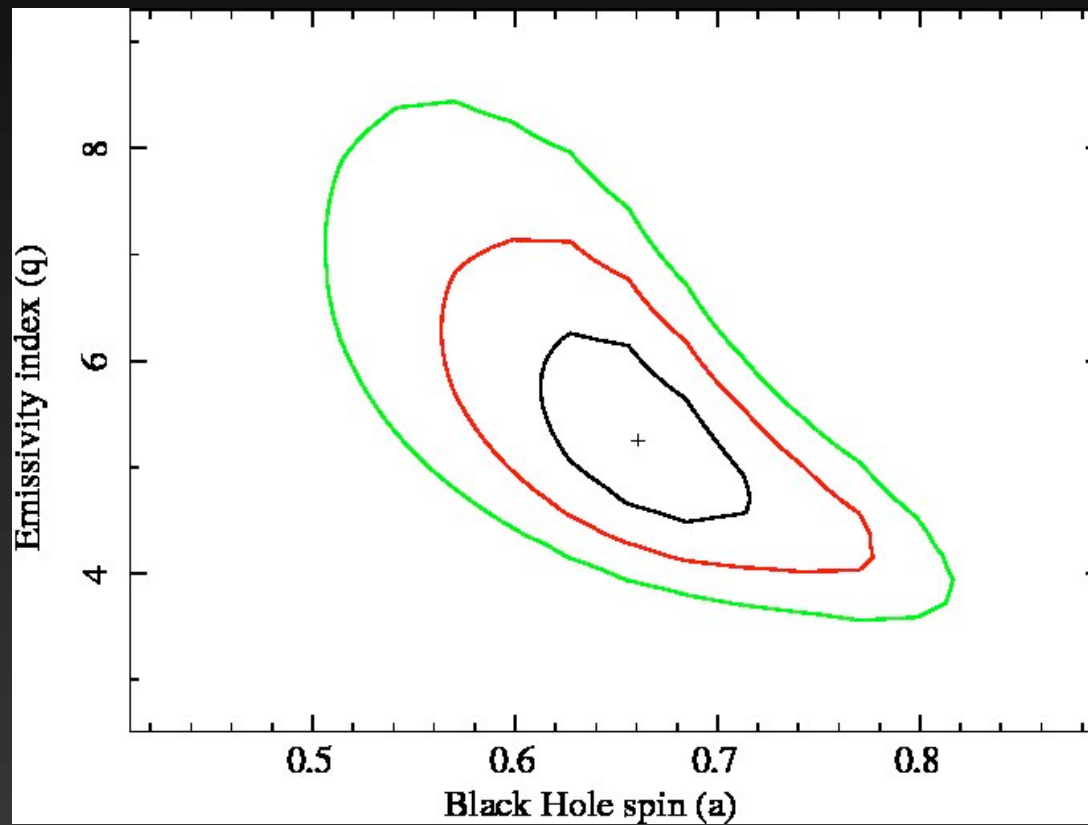
The hard X-ray flux ($\sim 2.5 \cdot 10^{-11}$ in 20-100 keV) makes it a good Suzaku source (especially for the HXD)





Swift J2127.4+5654 with Suzaku

modeling the XIS spectra with a relativistic line provides constraints on both **emissivity index** (steep) and **BH spin** (intermediate)



GM et al 09



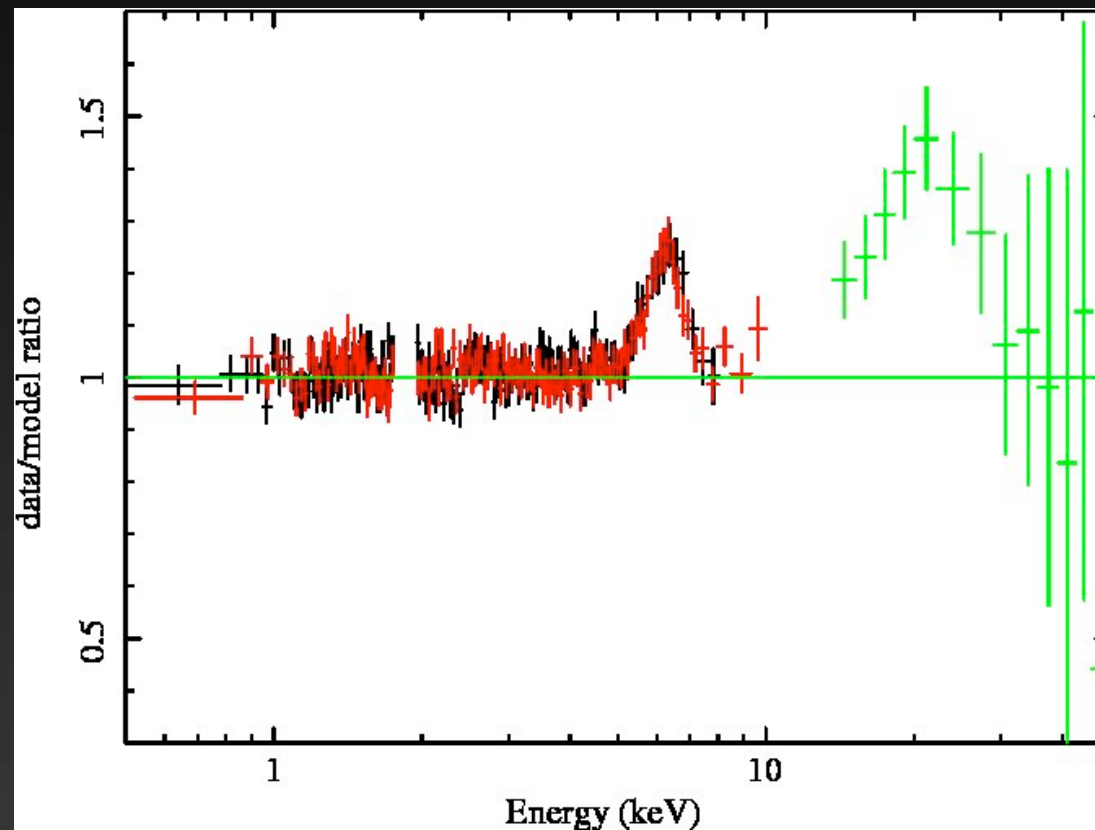
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The relativistic Fe K line should be associated with a full X-ray reflection spectrum which shows up at > 10 keV (Compton hump)



Swift J2127.4+5654 with Suzaku

The relativistic Fe K line should be associated with a **full X-ray reflection spectrum which shows up at > 10 keV** (Compton hump)

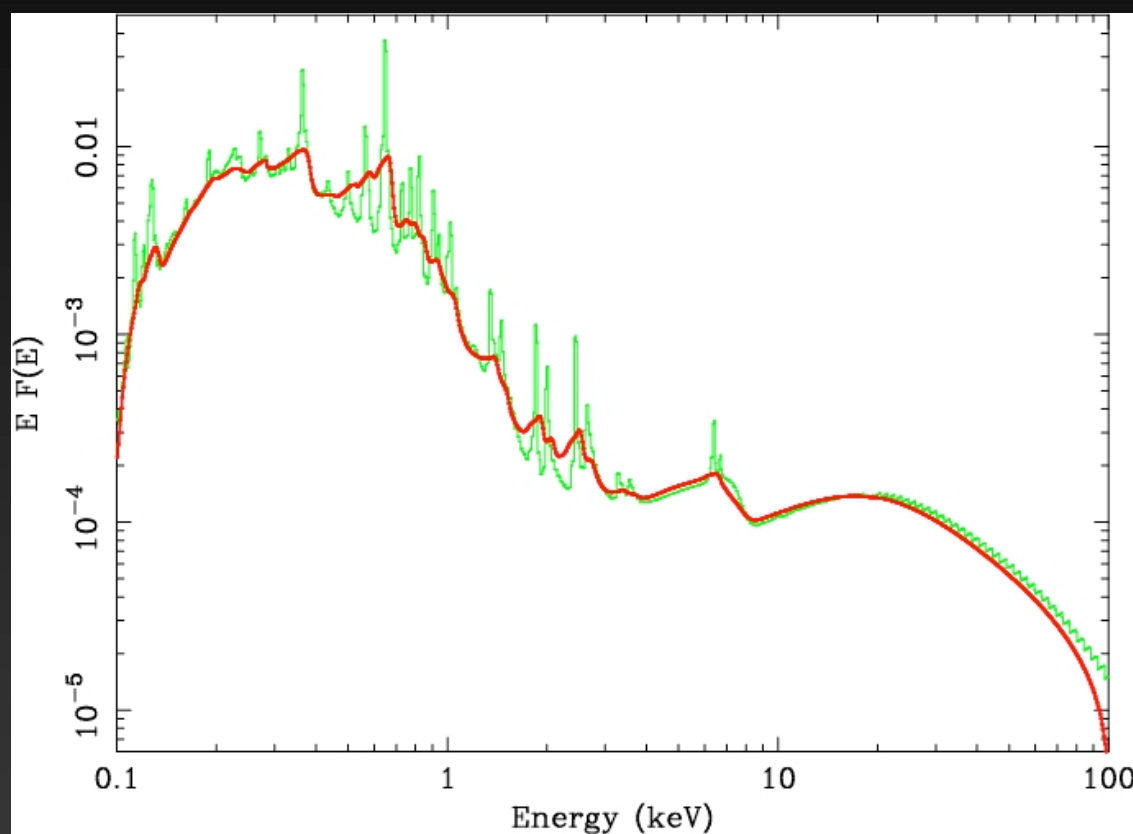


GM et al 09



Swift J2127.4+5654 with Suzaku

We then used a **ionized reflection model** (Ross & ACF 05) to account for the whole broadband spectrum to get the spin



Ross & Fabian 05

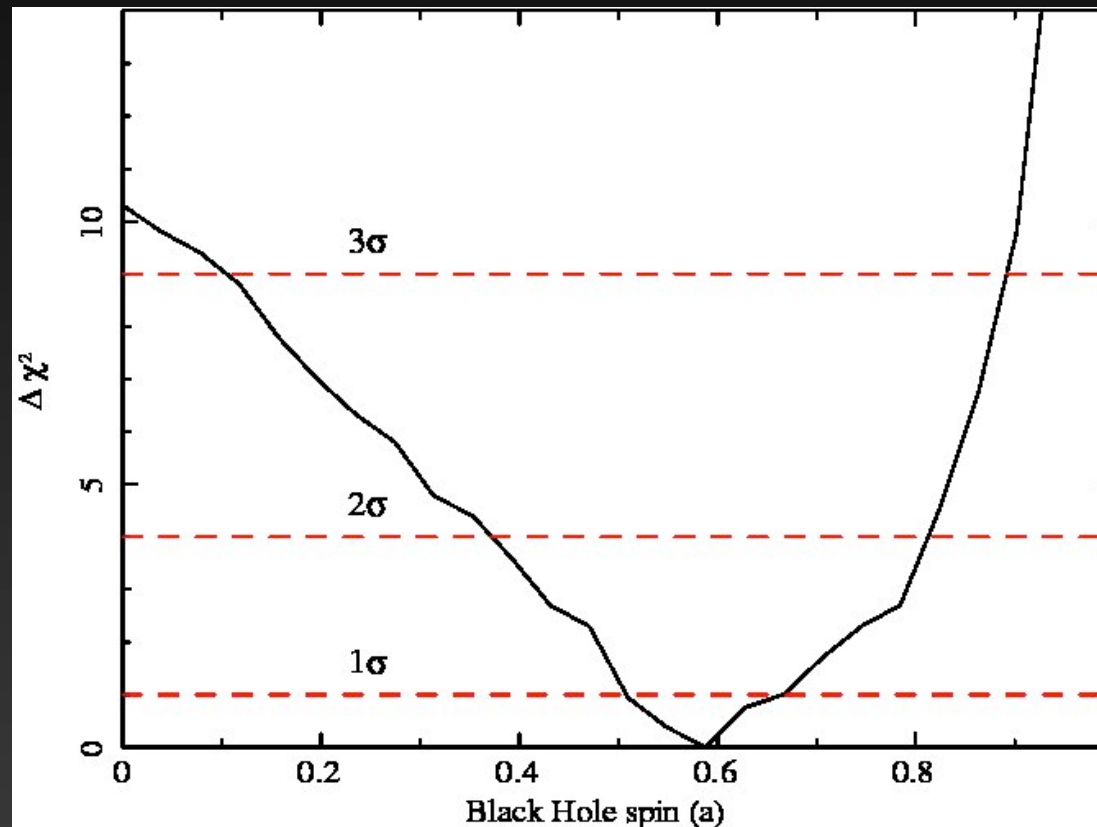


Swift J2127.4+5654 with Suzaku

The broadband analysis confirms results from Fe K diagnostics

$a \sim 0$ is excluded but just at the 3σ level

$a \sim 0.998$ is excluded at more than 5σ



GM et al 09



Swift J2127.4+5654 with Suzaku

BH accreting more than their own mass should be maximally spinning
this can then be excluded

a relatively **recent major merger** could result in an intermediate spin

$a \sim 0$ cannot be excluded with very high significance

if we also consider that the spin may be 20% overestimated

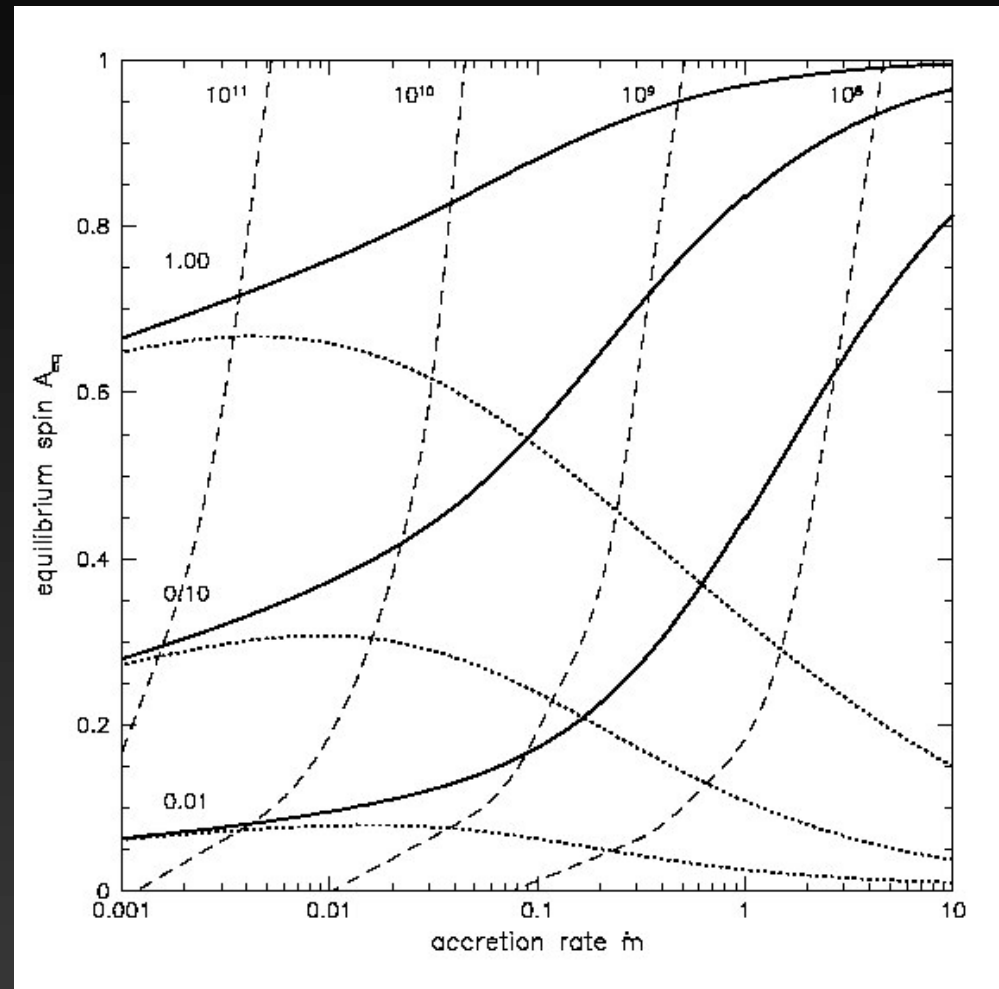
it could be that BH growth proceeds here **via chaotic accretion**

a further possibility: **magnetic extraction of rotational energy (BZ) ?**



Swift J2127.4+5654 with Suzaku

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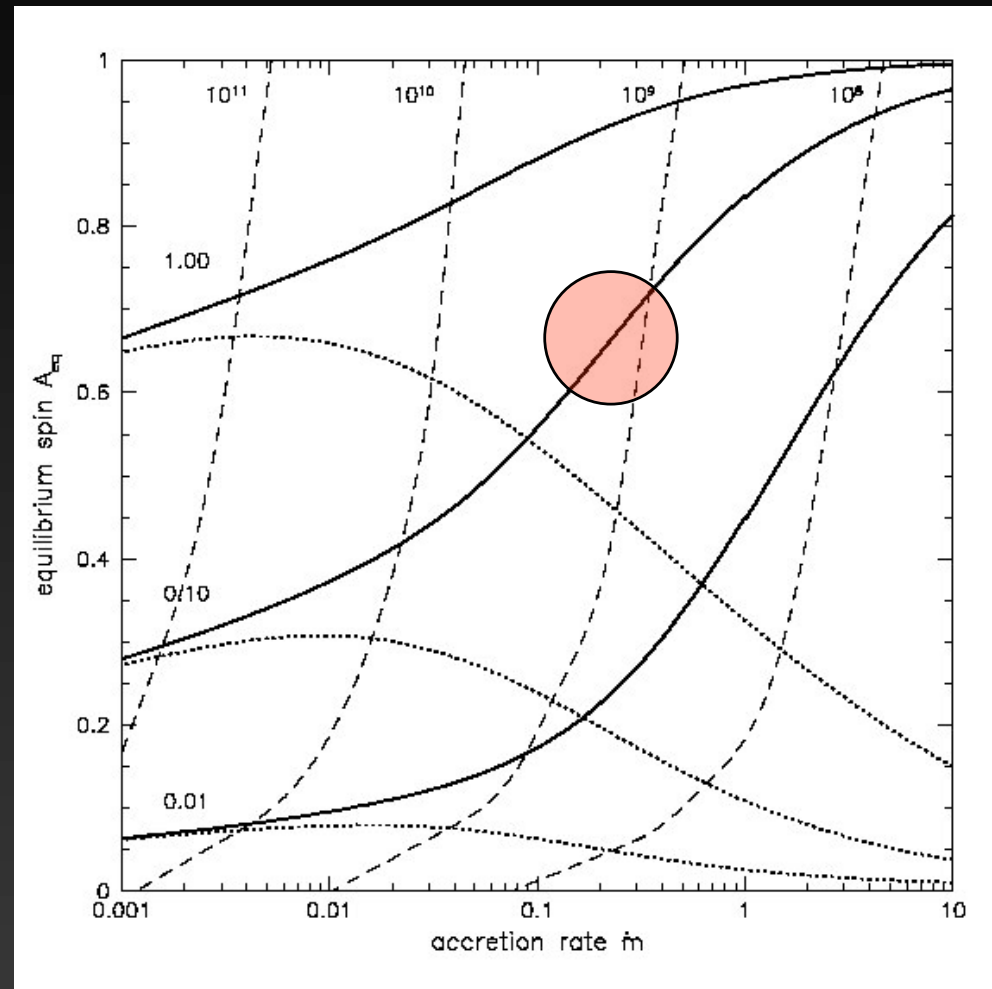


Moderski, Sikora & Lasota 98



Swift J2127.4+5654 with Suzaku

a further possibility: magnetic extraction of rotational energy (BZ) ?



Indeed, if we consider a standard $\alpha \sim 0.1$ disc coupled with the mass accretion rate of the object (~ 0.2) the system should reach an equilibrium value of $a \sim 0.5-0.6$, not far from what we get



Fairall 9 with Suzaku

Now a standard broad line Seyfert 1

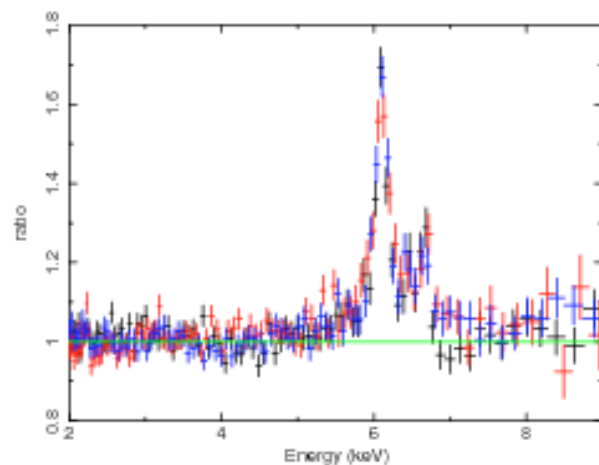
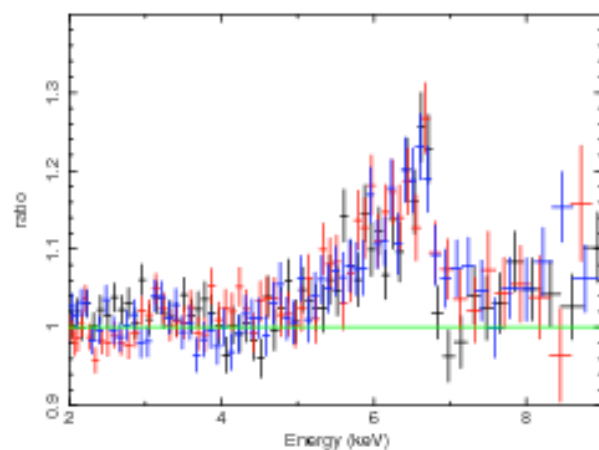


FIG. 2.— The plot above shows the data/model ratio in Fe K region that results from a simple power-law fit to the data. The narrow Gaussian peak near 6.1 keV (6.4 keV in the rest frame) is due to reflection from distant gas. A broad diskline component is also clearly present. The XIS0, XIS1, and XIS3 spectra are shown in black, red, and blue, respectively.





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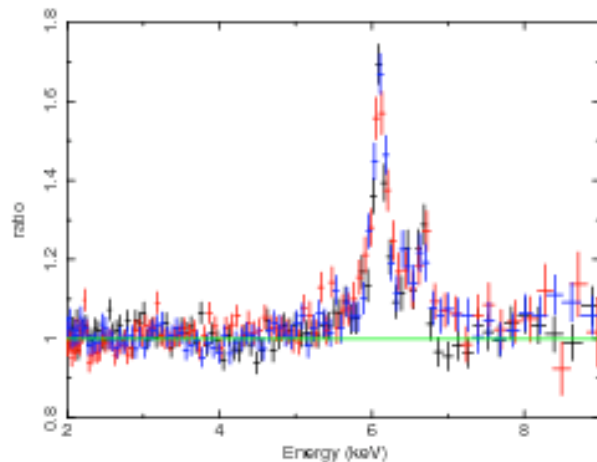
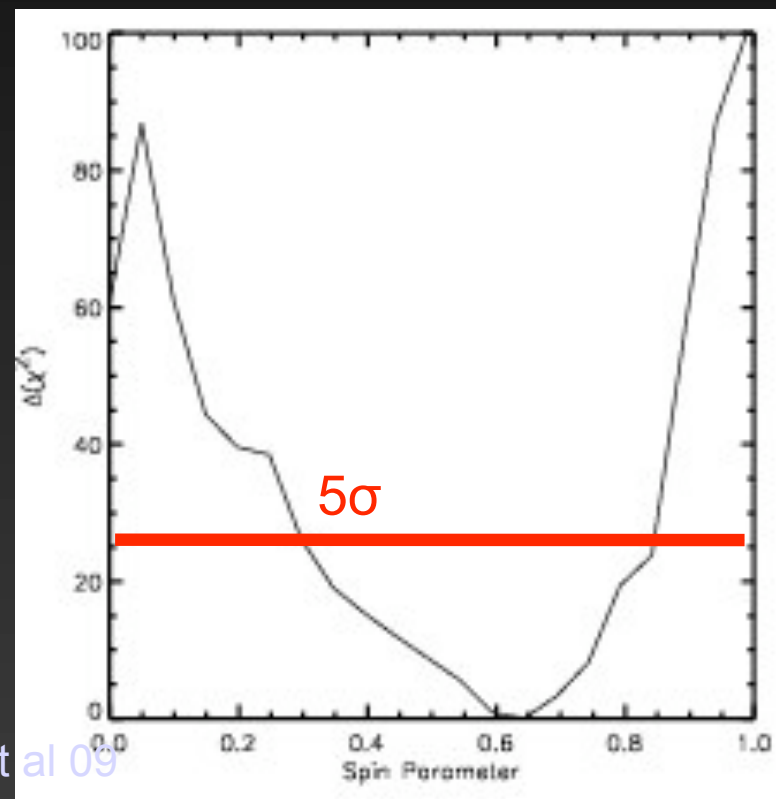
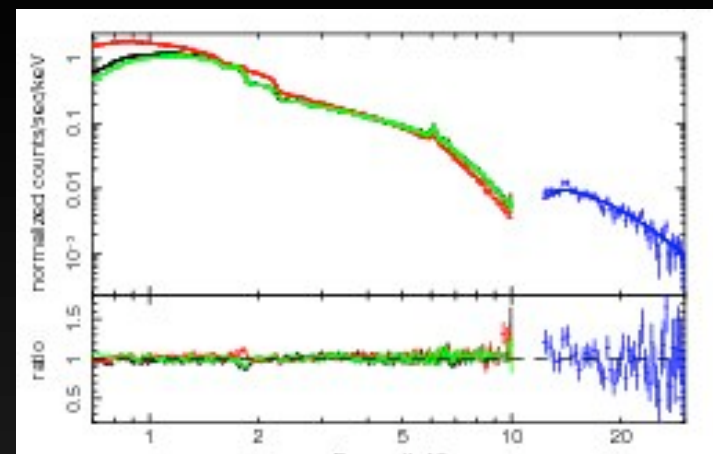
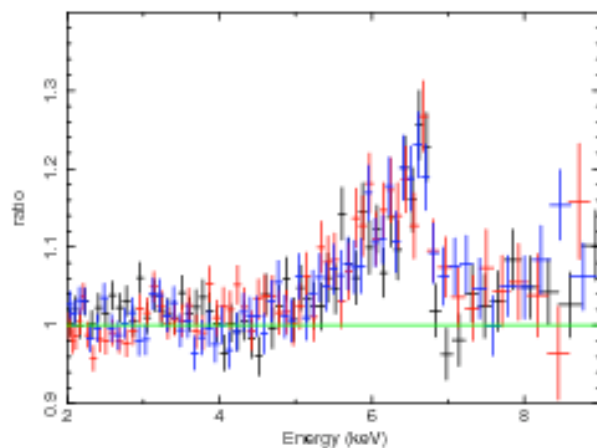


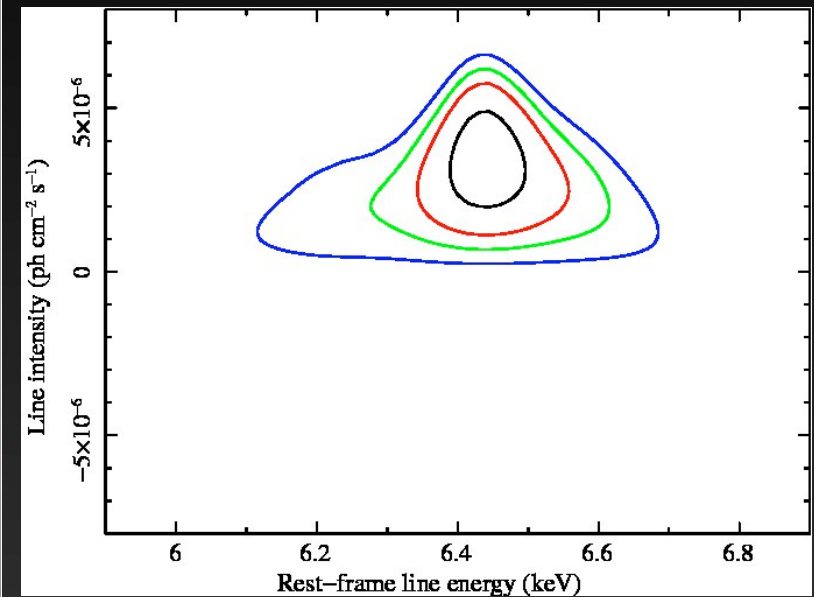
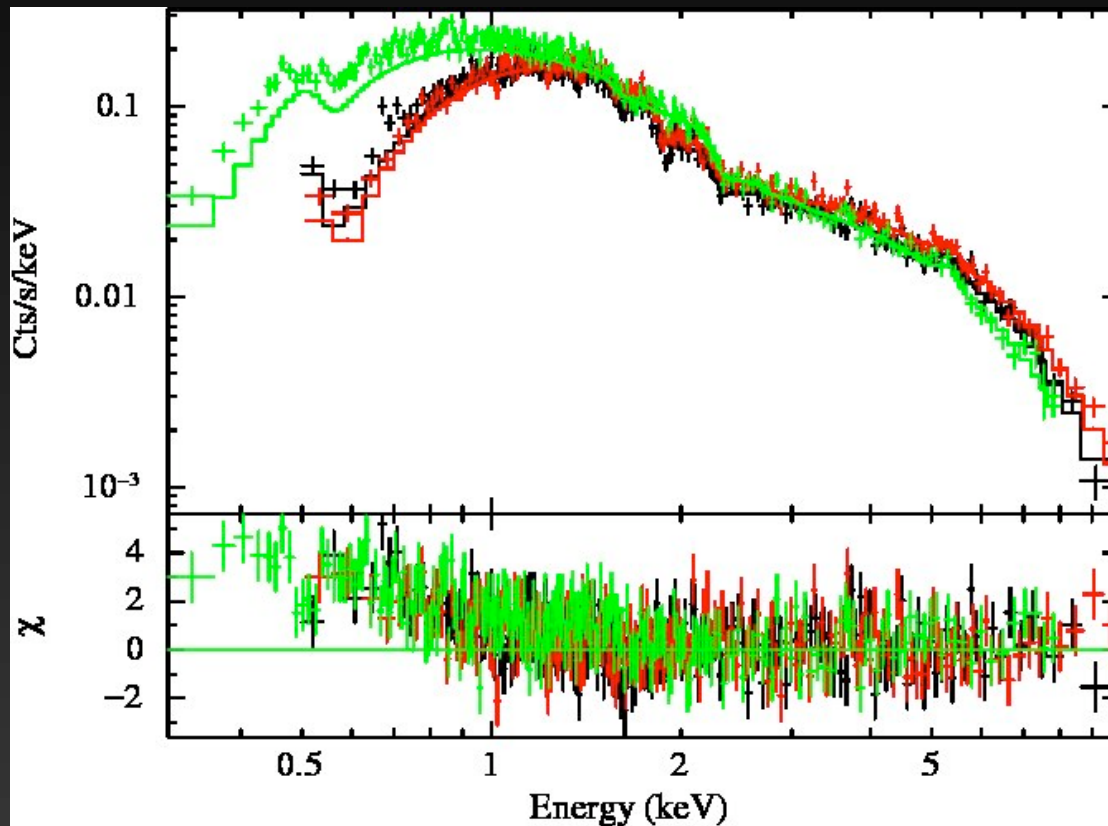
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BH spin from soft excess: RBS 1124

Now a broad line luminous QSO ($L / L_{\text{edd}} \sim 1$ from x-rays)



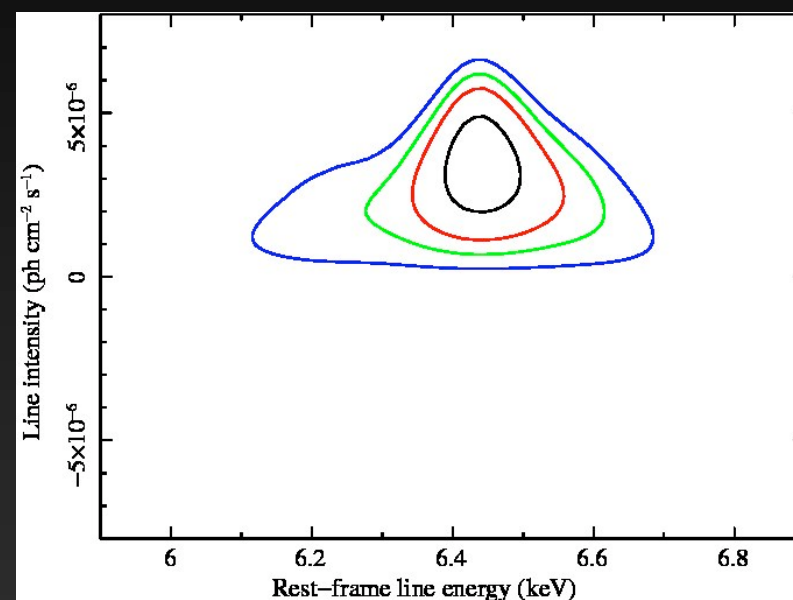
GM et al 10



BH spin from soft excess: RBS 1124

As far as the weak Fe K line is concerned, **no need for reflection from within $\sim 100 R_g$**

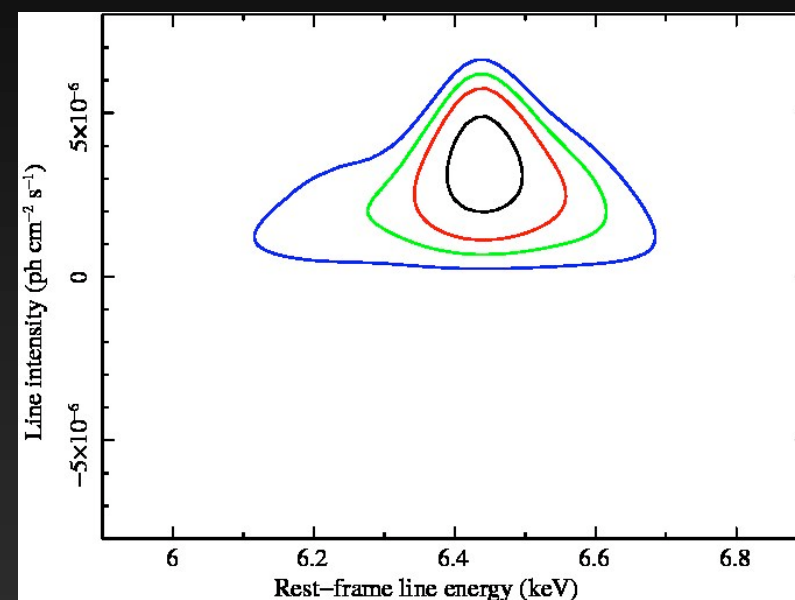
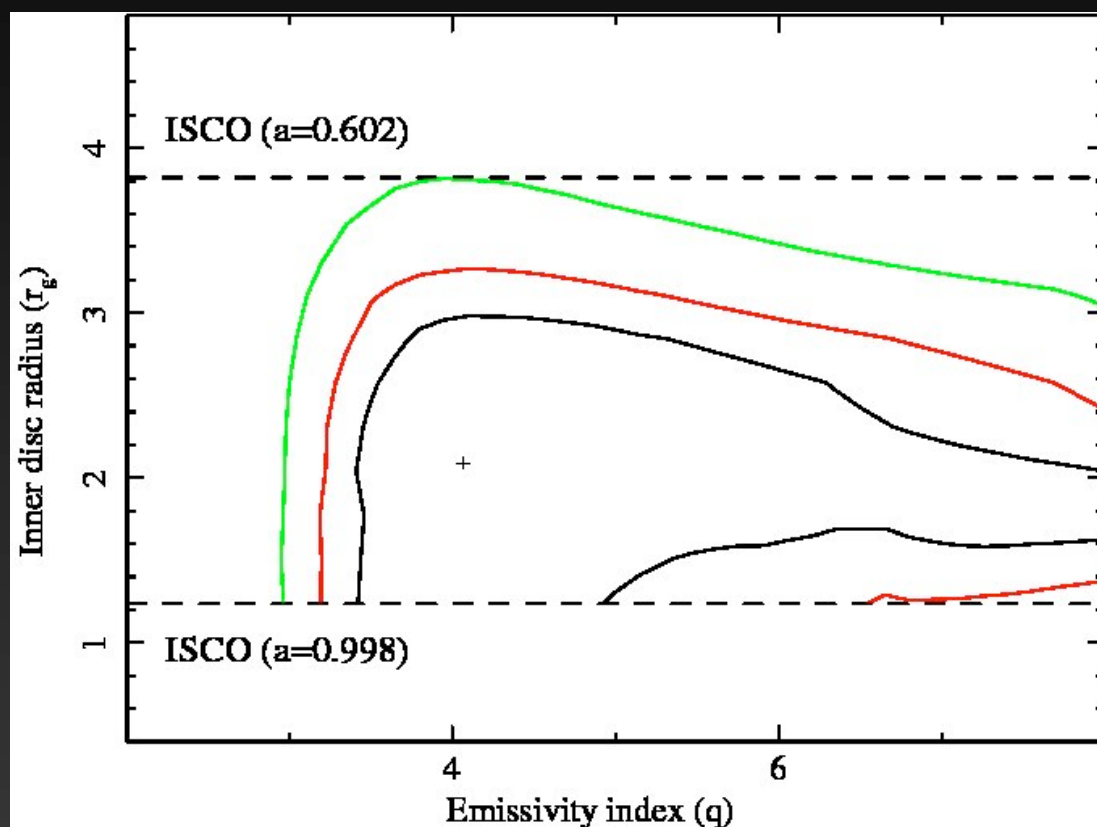
But **what about the soft excess ?**





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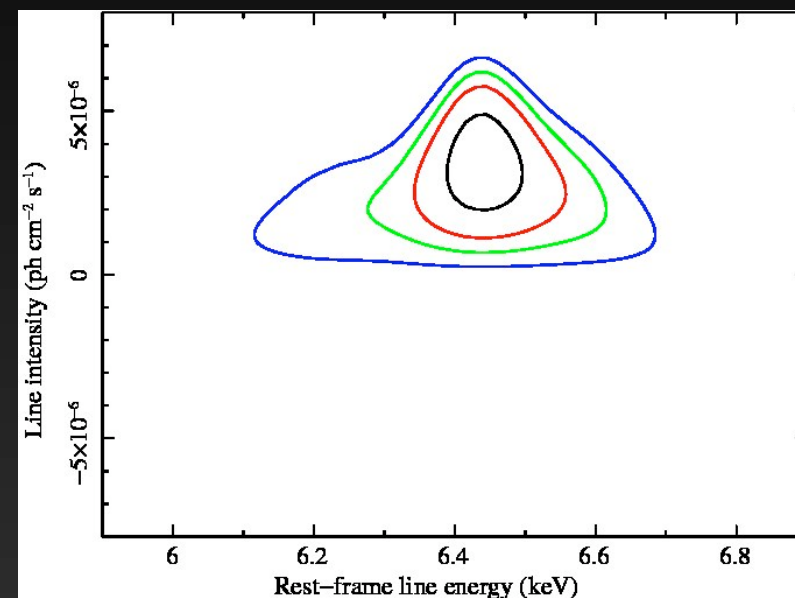
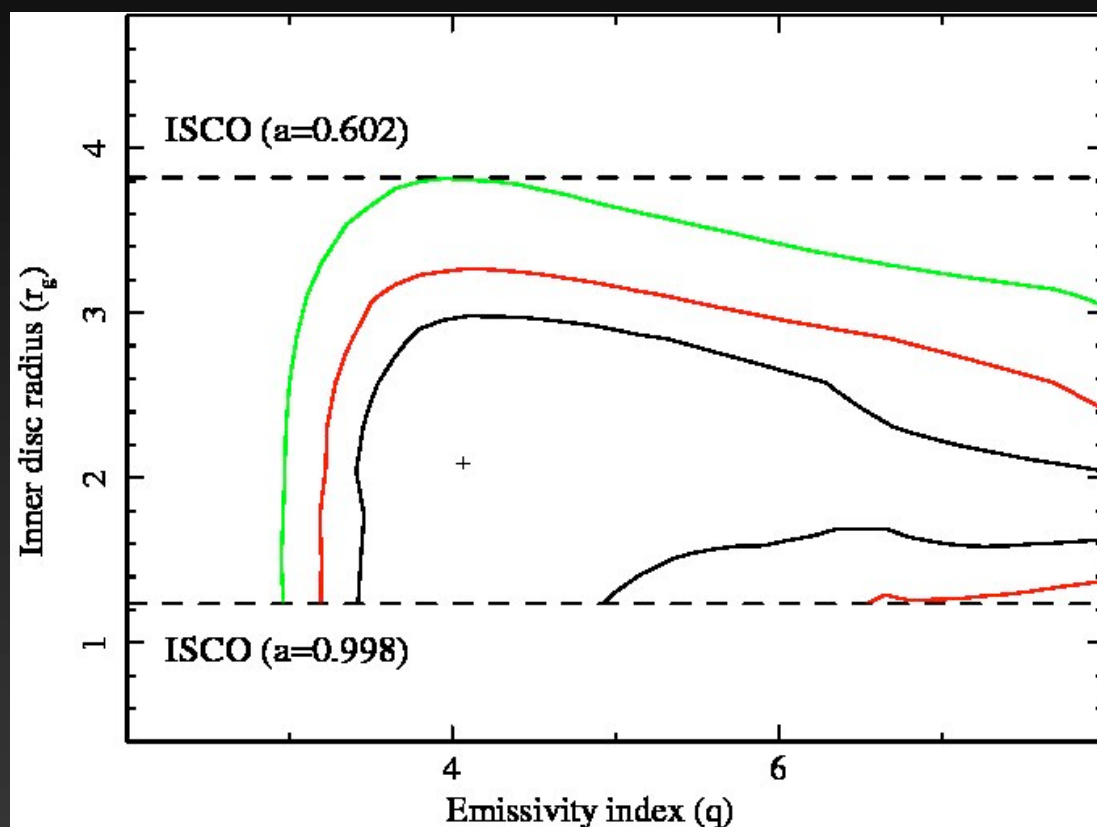




BH spin from soft excess: RBS 1124

Soft excess and Fe K are consistent, but still... how can it be ?

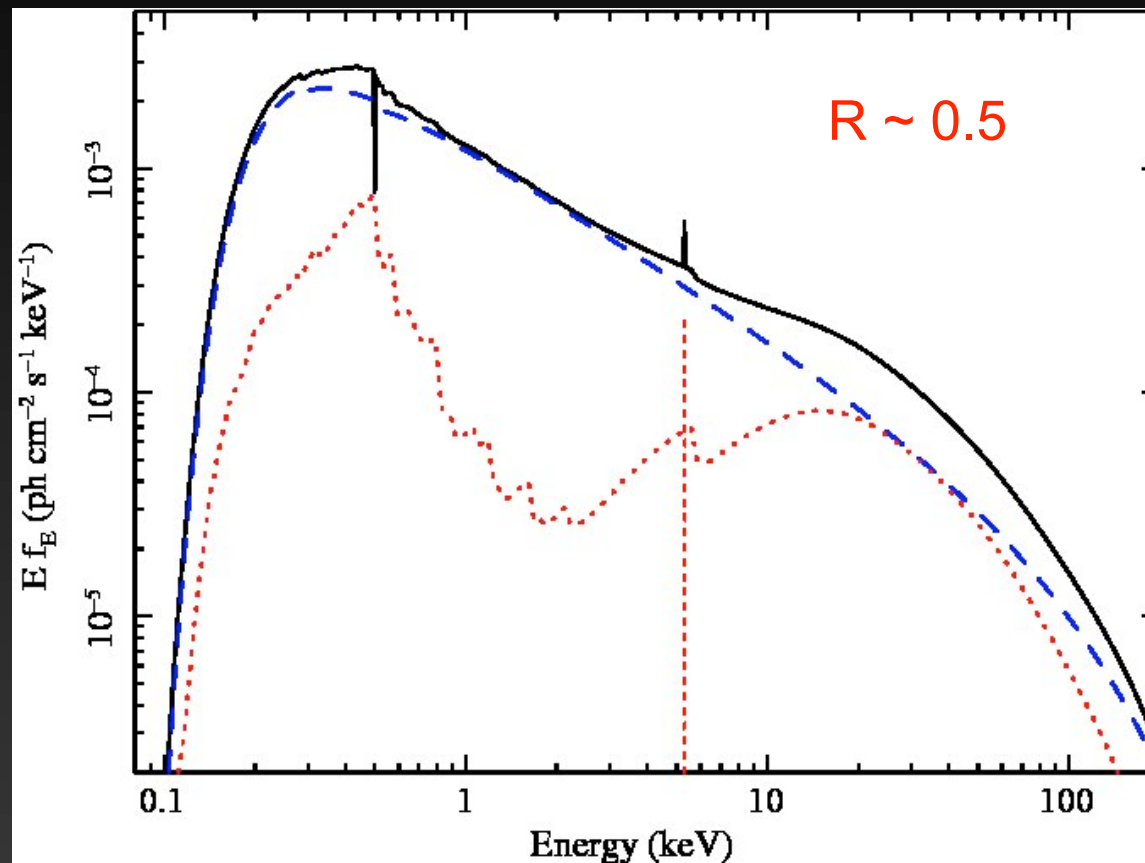
Remember, **the soft excess is weak in this object** -> **low R** ? If so, the Fe K we are able to detect is just the blue peak, red wing buried in continuum





BH spin from soft excess: RBS 1124

Not surprisingly, this is what we get from the modeling ($R \sim 0.5$)
Moreover $L/L_{\text{edd}} \sim 0.7$ (very efficient, disc should be at the ISCO)



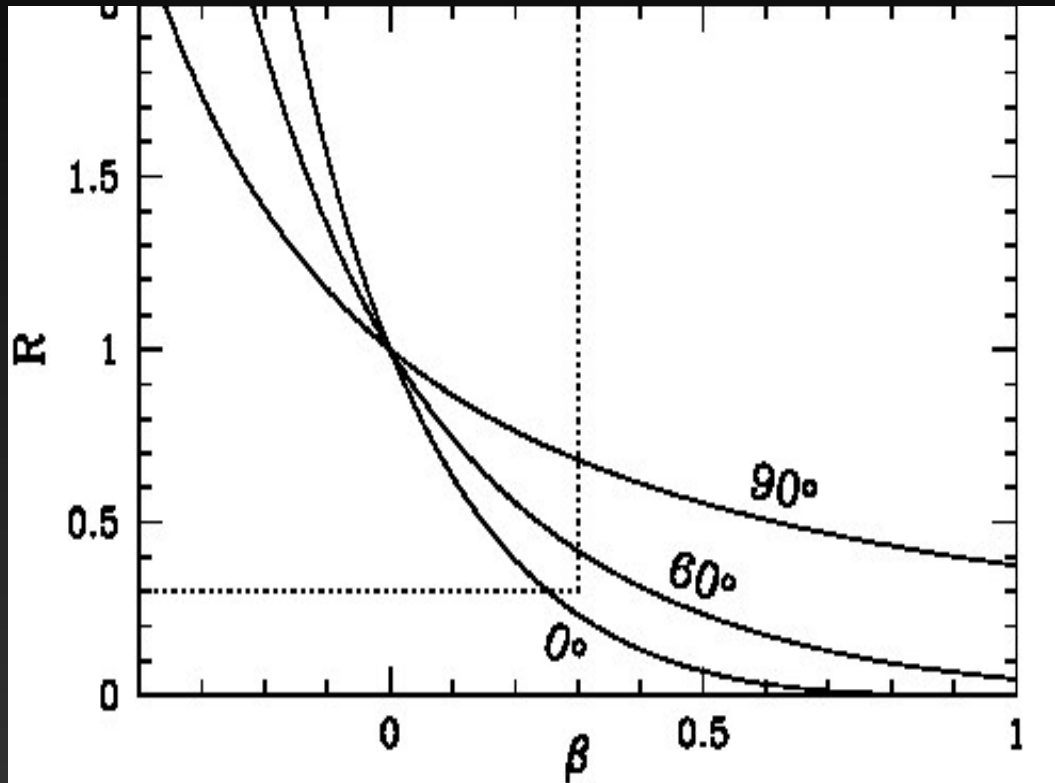
So, the right question to ask could be:

why such a low R ?



BH spin from soft excess: RBS 1124

The most natural answer (to me) is **mild X-ray beaming**



Beloborodov 99

Moreover this would reduce the intrinsic X-ray luminosity by a factor

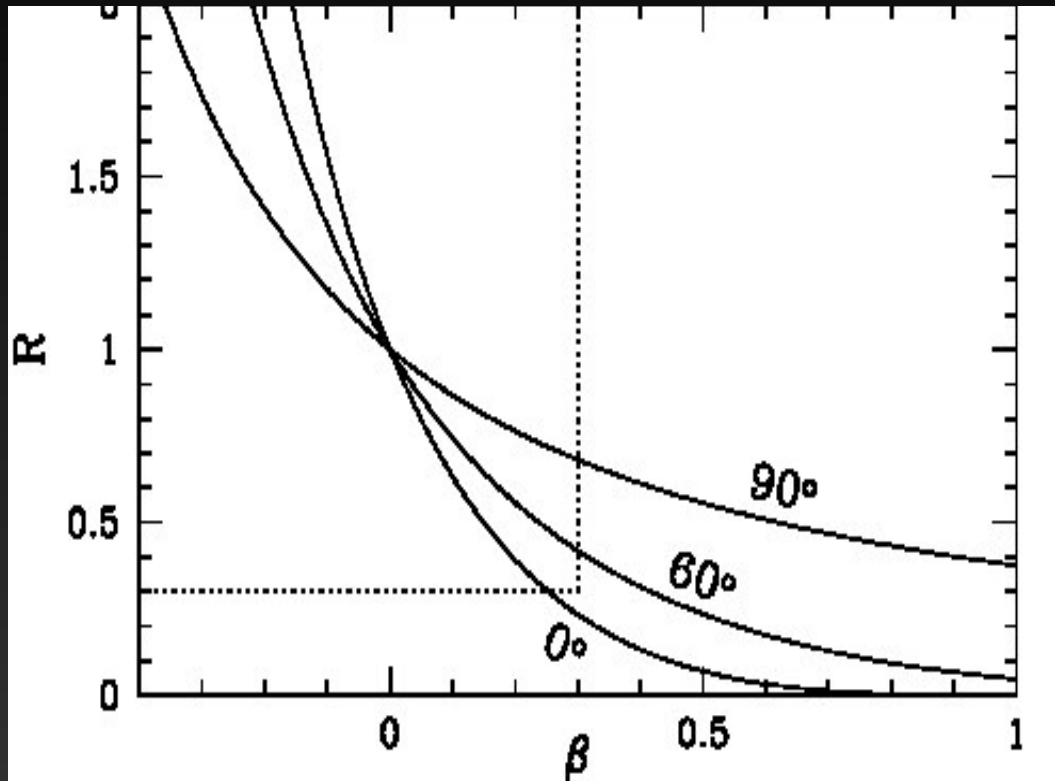
$$[\gamma (1 - \beta \cos i)]^{-3} \sim 0.5$$

leading to a more reasonable Eddington ratio ~ 0.4 , more consistent with optical luminosity



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Thanks !!